

Statistical Physics Summary for the Final Exam

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Part 1

Microcanonical ensemble (N,V,E)

1. Equal of priori probabilities:

Size of the phase space

$$\Omega(E) = \frac{1}{N! h^{3N} \int_{E < H < E + \Delta E} dq dp}$$

2. Observables:

$$\bar{A} = \frac{1}{\Omega} \frac{1}{N! h^{Nr} \int A(q, p) d\Omega}$$

3. Why ensembles?

4. $S = k \ln \Omega$

$$5. \begin{cases} d \ln \Omega = \beta dE + \gamma dV + \alpha dN \\ dU = TdS - PdV + \alpha N \\ lnm! = m lnm - m \end{cases}$$

Part 2

Canonical Ensemble (N,V,T).

$$1. \rho = \frac{1}{Z} e^{-\beta E_s}.$$

Considering the degeneration: $\rho_l = \frac{\omega_l}{Z} e^{-\beta E_l}$

$$2. \bar{A} = \frac{1}{Z} \sum_l A_l \omega_l e^{-\beta E_l} = \sum_l A_l \rho_l$$

$$3. \text{Tricks: } \sum_l \omega_l \rightarrow \sum_{q,p} \frac{\Delta p \Delta q}{N! h^{Nr}} \rightarrow \frac{1}{N! h^{Nr}} \int dp dq , \quad \text{here } N! \text{ comes from the}$$

undistinguishable particles.

$$4. \quad \left\{ \begin{array}{l} P = \frac{1}{\beta} \frac{\partial \ln Z}{\partial V} (\bar{Y} = -\frac{1}{\beta} \frac{\partial \ln Z}{\partial y}) \\ F = -\frac{1}{\beta} \ln Z \\ S = k(\ln Z - \beta \frac{\partial \ln Z}{\partial \beta}) \end{array} \right.$$

5. Microcanonical ensemble and Canonical ensemble are equivalent to each other at thermal dynamic limit.

6. Equipartition theorem. (For Classical systems.)

$$7. \quad \text{Virial theorem: } \left\langle q_i \frac{\partial H}{\partial q_i} \right\rangle = \langle q_i \dot{p}_i \rangle = kT .$$

8. Heat capacity:

a. Harmonic oscillators:

Write down the Hamiltonian;

Calculate the partition function Z;

Calculate the observables according to the following formulas:

$$\left\{ \begin{array}{l} U = -\frac{\partial \ln Z}{\partial \beta} \\ C_v = \frac{\partial U}{\partial T} \\ F = -\frac{1}{\beta} \ln Z \\ \mu = \frac{\partial F}{\partial N} \\ F = U - TS \rightarrow S = \textcolor{red}{i} \end{array} \right. ;$$

9. Heat capacity of solids:

Write down the decoupled Hamiltonian \rightarrow Cal. Z \rightarrow Cal. U.

Then the road diverges!

a. Einstein Model: $\omega_i = \omega$

b. Debye Model: (or continuous model) Assume the frequency distribution is:

$$g(\omega) d\omega = B\omega^2 d\omega, \text{ with } \omega_D \text{ as the cutoff}$$

$$(\int_0^{\omega_D} g(\omega) d\omega = 3N).$$

Part 3

Quantum statistics.

1. DM: $\hat{\rho} = \frac{1}{N} \sum_{k=1}^N |\Psi_k\rangle\langle\Psi_k|$, in which $a_k^n = \langle\phi_n|\Psi_k\rangle$.

2. Von Neumann Eq: $i\hbar\dot{\rho} = [\hat{H}, \hat{\rho}]$. $\dot{\rho} = 0$ when the system is under equilibrium.

3. Microcanonical ensemble: $\rho_n = \begin{cases} \frac{1}{\Gamma} & \\ \Gamma & \\ 0 & \end{cases}$

Canonical ensemble: $Z = e^{\beta\mu}$. $\rho = \frac{e^{-\beta H}}{\{tre^{-\beta H}\}}$. $Q(\mu, V, T) = \sum_N = 0^\infty z^N Q_N(\beta)$

$$\langle \hat{A} \rangle_c \equiv \text{tr}\{\hat{A}\hat{\rho}\}.$$

4. Details:

a. Diagonalise Hamiltonian.

b. Cal. DM: $\hat{\rho}$

c. Cal. Observables: $\langle \hat{A} \rangle_c$

5. Thermal length: $\lambda \equiv \left\{ \frac{\hbar^2}{2\pi mkT} \right\}^{1/2}$.

Part 4

1. Universality: deep rooted in the correlation length.
2. Critical exponent.
3. Ising Model:

a. Landau:

$$\left\{ \begin{array}{l} \mu(T, m) = \mu_0(T) + a(T)m^2 + \frac{1}{2}c(T)m^4 + \text{etc.} \\ a = a_1 t, c = c_0 \\ \text{Equilibrium } \frac{\partial \mu}{\partial m} = 0, \frac{\partial^2 \mu}{\partial m^2} > 0 \\ \Rightarrow m \\ \mu(T, m) = -Bm + \mu_0(T) + a(T)m^2 + \frac{1}{2}c(T)m^4 + \text{etc.} \\ \Rightarrow \chi = \frac{\partial m}{\partial B} \end{array} \right.$$

b. MFT:

$$(1) \text{ Hamiltonian: } H = -J \sum_{\langle i,j \rangle} \sigma_i \sigma_j - \mu B \sum_i \sigma_i$$

$$\Rightarrow H = \frac{1}{2} J N \bar{\sigma}^2 z - J z \bar{\sigma} \sum_i \sigma_i - \mu B \sum_i \sigma_i = \frac{1}{2} 2 J N z \bar{\sigma}^2 - \mu (B + \frac{J}{z} \mu \bar{\sigma}) \sum_i \sigma_i$$

$$(2) \text{ Partition func.: } Z_N = Z_i^N = \{e^{-\beta \mu B_{\text{eff}}} + e^{\beta \mu B_{\text{eff}}}\}^N$$

$$(3) \text{ Constrains: } \bar{\sigma} = \frac{1}{\beta} \frac{\partial}{\partial H} \ln Z$$

(4) Observables:

$$T_c, \text{ when } B=0 \rightarrow \bar{\sigma} = \tanh\left(\frac{zJ}{kT} \bar{\sigma}\right) \rightarrow T_c = \frac{zJ}{k}$$

$$\text{Near the critical zone, } \bar{\sigma} \leftarrow \tanh(\beta \mu B_{\text{eff}})$$

$$C(B=0) = \frac{dU}{dT}, \quad U = \langle H \rangle = -\frac{1}{2} J z \bar{\sigma}^2 N$$

$$\chi(B, T) = \frac{\partial \bar{M}}{\partial B}_T = N \mu \left(\frac{\partial \bar{\sigma}}{\partial B} \right)_T$$

c. Transfer Matrix method.

d. Recursion method. $Z_{N+1} = f(Z_N)$

e. Renormalization method.

4. Order parameter.